On the sup-norm of Maass cusp forms of large level GERGELY HARCOS (joint work with Nicolas Templier)

Comparing the various Sobolev norms of automorphic forms is useful in the theory of quantum chaos and subconvexity of L-functions, which in turn have deep arithmetic applications. We consider the following special case.

Problem. Let f be a Hecke–Maass cuspidal newform of level N and Laplacian eigenvalue λ . Assume that $||f||_2 = 1$ with respect to $dxdy/y^2$. Bound $||f||_{\infty}$ in terms of N and λ .

In the λ -aspect the first nontrivial (and so far unsurpassed) bound is due to Iwaniec and Sarnak [6]: $||f||_{\infty} \ll_{N,\epsilon} \lambda^{5/24+\epsilon}$ for any $\epsilon > 0$. In the *N*-aspect the trivial bound is $||f||_{\infty} \ll_{\lambda,\epsilon} N^{\epsilon}$, while the most optimistic bound would be $||f||_{\infty} \ll_{\lambda,\epsilon} N^{-1/2+\epsilon}$. Here and later, the dependence on λ is understood to be continuous. The breakthrough in the *N*-aspect was recently achieved by Blomer– Holowinsky [2] who proved $||f||_{\infty} \ll_{\lambda,\epsilon} N^{-25/914+\epsilon}$, at least for square-free *N*. The restriction on *N* seems difficult to remove as it is needed for a certain application of Atkin–Lehner theory. By a systematic use of geometric arguments Templier [7] derived $||f||_{\infty} \ll_{\lambda,\epsilon} N^{-1/22+\epsilon}$, and Helfgott–Ricotta [3] improved this to $||f||_{\infty} \ll_{\lambda,\epsilon} N^{-1/20+\epsilon}$. As we shall explain below, an efficient use of Atkin– Lehner theory leads to a short and clean proof of the following result [5]:

Theorem. Let f be an L^2 -normalized Hecke–Maass cuspidal newform of squarefree level N, trivial nebentypus, and Laplacian eigenvalue λ . Then for any $\epsilon > 0$ we have a bound

$$\|f\|_{\infty} \ll_{\lambda,\epsilon} N^{-1/6+\epsilon},$$

where the implied constant depends continuously on λ .

The theorem improves our earlier bound [4] with exponent $-1/12 + \epsilon$. A hybrid version can also be established, improving significantly on [2, Theorem 2].

We turn to an informal discussion of our method. Very vaguely, the idea of proving a result as above has been like this:

- (1) Pick any $z \in \mathfrak{H}$ where |f(z)| needs to be estimated.
- (2) Apply an Atkin–Lehner operator on z to ensure that Im z is not too small.
- (3) Use the amplification method and some trace formula to reduce the problem to a counting problem depending on z.
- (4) Do the counting based on the diophantine properties of z.

Our improvement results mainly from the following shortcut:

- (2) Apply an Atkin–Lehner operator on z to maximize Im z.
- (4') Observe that z has good diophantine properties automatically, allowing a more efficient counting.

For a square-free level N the Atkin–Lehner operators can be represented by matrices of the form

$$W_M = \frac{1}{\sqrt{M}} \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \mathrm{SL}_2(\mathbb{R}), \qquad M \mid N,$$

where $a, b, c, d \in \mathbb{Z}$ are integers satisfying

$$ad - bc = M,$$
 $a \equiv 0 (M),$ $d \equiv 0 (M),$ $c \equiv 0 (N).$

A key feature is the multiplication rule

$$W_M W_{M'} = W_{M''}$$
 with $M'' = \frac{MM'}{(M,M')^2}$

which shows that the W_M 's form a group $A_0(N)$ containing $\Gamma_0(N)$ as a normal subgroup. As a result, Atkin–Lehner operators induce an action on $\Gamma_0(N) \setminus \mathfrak{H}$ by the finite group $A_0(N)/\Gamma_0(N) \cong (\mathbb{Z}/2\mathbb{Z})^{\omega(N)}$, where $\omega(N)$ is the number of distinct prime factors of N.

By Atkin–Lehner theory [1], a Hecke–Maass cuspidal newform f of level N is an eigenvector for $A_0(N)$ with eigenvalues ± 1 , therefore in examining the sup-norm of f we can restrict to the following fundamental domain for $A_0(N)$:

$$\mathcal{F}(N) := \{ z \in \mathfrak{H} \mid \operatorname{Im} z \ge \operatorname{Im} \delta z \text{ for all } \delta \in A_0(N) \}.$$

Our starting point was the observation that the elements of $\mathcal{F}(N)$ have good diophantine properties (we assume that N is square-free):

Lemma. Let $z = x + iy \in \mathcal{F}(N)$. Then the lattice $\langle 1, z \rangle$ has minimal distance at least $N^{-1/2}$ and covolume $y \gg N^{-1}$.

The usefulness of this lemma becomes apparent when we relate |f(z)| to a lattice counting problem depending on z. By combining the amplification method of Duke–Friedlander–Iwaniec with the pretrace formula of Selberg we obtain

$$\Lambda^2 |f(z)|^2 \ll_{\lambda,\epsilon} N^{\epsilon} \sum_{l \ge 1} \frac{y_l}{\sqrt{l}} M(z,l,N),$$

where M(z, l, N) denotes the number of integral matrices $\gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$ such that

(*)
$$\det(\gamma) = l, \quad c \equiv 0 (N), \quad \left| -cz^2 + (a-d)z + b \right|^2 \leq ly^2 N^{\epsilon},$$

 Λ is a large parameter (the amplifier length), and

$$y_l := \begin{cases} \Lambda, & l = 1; \\ 1, & l = l_2 \text{ or } l_1 l_2 \text{ or } l_1 l_2^2 \text{ or } l_1^2 l_2^2 \text{ with } \Lambda < l_1, l_2 < 2\Lambda \text{ primes}; \\ 0, & \text{otherwise.} \end{cases}$$

Our second key observation is that for each range $l \simeq L$ and for each fixed c the inequality in (*) can be used to bound the number of choices for the pair (a-d,b).

Indeed, the pairs correspond to the lattice points in a disk of radius $\ll L^{1/2}yN^{\epsilon}$, hence the Lemma together with some geometry of numbers yields the bound

$$#(a-d,b) \ll_{\epsilon} N^{\epsilon} \left(1 + L^{1/2} N^{1/2} y + L y\right)$$

for any c. A simple manipulation of (*) also shows $c \ll_{\epsilon} L^{1/2} N^{\epsilon}/y$. Finally for any triple (c, a - d, b) we regard the identity

$$(a+d)^2 - 4l = (a-d)^2 + 4bc$$

as an equation for the pair (a+d, l). By the sparsity of potential *l*'s we can bound the number of quadruples (c, a-d, b, a+d) efficiently, which of course is the same as bounding the sum of M(z, l, N) over the *l*'s considered.

Along these lines we obtain

$$\Lambda^2 |f(z)|^2 \ll_{\lambda,\epsilon} N^{\epsilon} \left(\Lambda + \Lambda^{5/2} N^{-1/2} + \Lambda^4 N^{-1} \right),$$

at least when $y < N^{-2/3}$ and $\Lambda^4 < y^{-2}N^{-\epsilon}$. The latter is automatic for $y < N^{-2/3}$ under the choice

$$\Lambda := N^{1/3 - \epsilon}$$

which incidentally also balances the terms in the previous display. Hence by amplification we really see that

$$f(x+iy) \ll_{\lambda,\epsilon} N^{-1/6+\epsilon}, \qquad y < N^{-2/3}$$

For the remaining range $y \ge N^{-2/3}$ we use a simple bound based on the Fourier expansion at the cusp ∞ (see [7, § 3.2] or [2, (92) & (27)]):

$$f(x+iy) \ll_{\lambda,\epsilon} N^{-1/2+\epsilon} y^{-1/2} \leqslant N^{-1/6+\epsilon}, \qquad y \geqslant N^{-2/3}.$$

References

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